### P3317 HW from Lecture 16+17 and Recitation 9

### Due Oct 30, 2018

# Problem 1. Born-Openheimer potential for $H_2^+$

The  $H_2^+$  ion consists of two protons, and one electron. Within the Born Oppenheimer approximation, we "integrate out" the electron to produce an effective potential felt by the protons. The Born-Oppenheimer potential is

$$V(R_1 - R_2) = \frac{e^2}{4\pi\epsilon_0} \frac{1}{|R_1 - R_2|} + \mathcal{E}(R_1 - R_2)$$
(1)

where the first term is the direct interaction between the protons, and the second term is the effective interaction from the electrons. In the case of  $H_2^+$ , this electronic contribution is the eigenvalue of a single electron problem,

$$\mathcal{E}\psi(r) = -\frac{\hbar^2}{2m_e} \nabla^2 \psi(r) - \frac{e^2}{4\pi\epsilon_0} \left[ \frac{1}{|R_1 - r|} + \frac{1}{|R_2 - r|} \right] \psi(r).$$
(2)

Adimensionalized, this equation reads

$$\mathcal{E}\psi(r) = -\frac{1}{2}\nabla^2\psi(r) - \left[\frac{1}{|R_1 - r|} + \frac{1}{|R_2 - r|}\right]\psi(r).$$
(3)

We will work with this adimensionalized form

**1.1.** How is the dimensionless energy related to the dimensionful energy? (Write the conversion factor in terms of the Rydberg energy  $E_R$ . Note: Wikipedia uses the symbol Ry instead of  $E_R$ .) How is the dimensionless length related to the dimensionful length? (Write the conversion factor in terms of the Bohr radius  $a_0$ .)

If you are interested in an analytic variational solution to Eq. (2) – you can look at Griffiths Chapter 7.3. It is kind of ugly though, because it is a 3D problem without spherical symmetry. A few years ago I had a homework problem that worked through it, but many students found it dissatisfying and mathematical. Instead we will numerically consider the 1D analog

$$\mathcal{E}\psi(x) = -\frac{1}{2}\partial_x^2\psi(x) - \left[W(x - X_1) + W(x - X_2)\right]\psi(x),$$
(4)

where W(x) is the Coulomb interaction between the nucleus and the electron. It turns out that the 1/x potential is not that well behaved in 1D, so we will use a truncated version

$$W(x) = \frac{1}{\sqrt{x^2 + \delta^2}}.$$
(5)

We will take  $\delta$  to be a small number, so that this is similar to a 1/x potential. For your numerics take  $\delta = 0.1$ .

**1.2.** Plot the function  $V_e(x) = -W(x-X_1) - W(x-X_2)$  for x running from -5 to 5 with  $X_1 = -3/2$  and  $X_2 = 3/2$ . Choose a scale for your axes so that the main features can be seen. Label your axes. Do not submit a program which plots this, just submit this graph.

**1.3.** Write a program which will numerically solve Eq. (4), finding the 5 eigenvalues closest to -8, for fixed  $d = X_1 - X_2$ . [These will be the 5 lowest eigenvalues.] Take  $\delta = 0.1$ , and a real space grid running from x = -10 to 10 with a grid spacing of 0.1. Place the two "protons" symmetrically,  $X_1 = d/2$  and  $X_2 = -d/2$ .

To make eigsh return eigenvalues close to -8, make one of its arguments sigma=-8.

Tabulate the function  $\mathcal{E}(d)$ , for these 5 states, with d running from 0 to 5 in steps of 0.1. Do not hand this in – this is just so that you can plot V.

Make a properly labeled plot of the Born-Oppenheimer potential  $V(d) = \mathcal{E}(d) + W(d)$  as a function of d for these 5 orbitals. [We add W(d), as this is the direct interaction between the protons.] Do not submit your program, just submit this graph with your homework. Annotate your graph (either by hand or with the computer) to point out the potential coming from the bonding orbital and from the lowest antibonding orbital.

**1.4.** Plot the electronic wavefunction for the lowest bonding and antibonding orbitals when the nuclei are separated by a distance d = 3. Label your graph, and mark which is bonding, and which is antibonding. *Do not submit your program, just submit the graph.* 

**Problem 2. Semiclassical Molecule** I want you to model the excitation spectrum of a neutral diatomic molecule. Generically there are three types of excitations:

electronic – an electron is promoted to a higher level

vibrational - the distance between the two atoms oscillates

rotational – the atoms rotate around one-another

Here you will estimate the energy scale of each of these excitations.

**2.1.** What is the typical bond length in a diatomic molecule (such as CO, or  $H_2$ )? Feel free to look this up in a textbook or on the web. You could also figure it out with dimensional analysis. We just need an order of magnitude.

**2.2.** One can roughly model the electronic states by considering the electrons to be trapped in a square box of size given by your solution to 2.1. If so, what is the energy scale of electronic excitations (in eV)?

One can estimate the vibrational energy by approximating the Born-Openheimer potential as a harmonic oscillator  $V(r) \sim \frac{1}{2}k(r-r_0)^2$ , where  $r_0$  is the equilibrium size of the molecule, and r is the separation between the nucleii. We will use dimensional analysis to estimate the constant k.

**2.3.** What units does k have?

2.4. The Born Oppenheimer potential comes from the solution of an electronic problem in the presence of the nuclei. List the dimensional quantities that it seems reasonable for k to depend on. Can you construct a unique k from these? If you have too many dimensional quantities, check to see if any of them are similar in size. Since we are looking for a numerical answer, these all are equivalent for us. If you still have too many, critically evaluate them and discard ones which are less likely to play a role in this physics. If you have too few dimensional quantities, look at Eq. (2). This is the equation for the Born-Oppenheimer potential, so they should all be there. Maybe you missed a fundamental constant like  $\hbar$ ?

Use these numbers to estimate k (in physical units).

**2.5.** The energy of these vibrations should be  $\hbar\omega = \hbar\sqrt{k/m}$ , where *m* should be a the nuclear mass. Estimate the energy of vibrational excitations (in eV).

**2.6.** For the rotational energy scales, we can use the classical expression,  $E = L^2/2I$ . Quantum mechanically, what are the allowed values of L?

2.7. From the geometry of the diatomic molecule, estimate the moment of inertia I.

**2.8.** Estimate the energy (in eV) of the rotational excitations.

**Problem 3. Paschen-Back effect** The electronic and nuclear spins of Rubidium-87 are coupled by the "hyperfine interaction." This is basically the dipole-dipole coupling between these spins. Of course, the typical magnetic moment of a nucleus is very very small, so the coupling is quite weak. Nonetheless, one can detect this splitting by "Magnetic Resonance" – the same technology which is used in the "Magnetic Resonance Imaging" used for medical diagnosis.

The model used for the Hyperfine splitting is

$$H = -g_e \mu_e \hat{S}_z B - g_{\rm Rb} \mu_N \hat{I}_z B + A_{hf} \hat{\mathbf{S}} \cdot \hat{\mathbf{I}}$$
(6)

where *B* is a magnetic field in the  $\hat{z}$  direction. Numerically  $\mu_e = 10^{-23}$  J/T, and  $\mu_N = 5 \times 10^{-27}$  J/T,  $g_e \approx 2/\hbar$ ,  $g_{\rm Rb} \approx 3/\hbar$ , and  $A_{hf} = \hbar \times 1.5$  GHz/ $\hbar^2$ . The nuclear spin is 3/2. and the electronic spin is 1/2. The operator  $\hat{\mathbf{S}} = (\hat{S}_x, \hat{S}_y, \hat{S}_z)$  corresponds to the electronic spin, and  $\hat{\mathbf{I}} = (\hat{I}_x, \hat{I}_y, \hat{I}_z)$  corresponds to the nuclear spin.

We want to understand the behavior of the energy levels as a function of the magnetic field.

**3.1.** Give two independent dimensionless numbers that can be made from  $\mu_e, \mu_N, B, A_{hf}, \hbar$ ? [Hint there are only two such numbers, and as usual they are not unique.]

**3.2.** Given that the ratio  $\mu_b/\mu_e$  is so small, it is reasonable to drop it from the problem, and consider the simplified model

$$H = -g_e \mu_e S_z B + A_{hf} S \cdot I. \tag{7}$$

Convert this equation to dimensionless units. Measure angular momenta in units of  $\hbar$ , energy in units of  $A_{hf}\hbar^2$  and B in units of  $A_{hf}\hbar/g_e\mu_e$ . Use this dimensionless expression for the rest of the problem. I do not want to see you write  $A_{hf}$  or g or  $\hbar$  for the rest of this problem. [Except for your graph at the end.]

For the following two subproblems it will be helpful to use the fact that for a spin-1/2 S the states obey: (note in our dimensionless units, all of these hbars are equal to unity)

$$S_z|S = 1/2, m_s = 1/2\rangle = (\hbar/2)|S = 1/2, m_s = 1/2\rangle$$
 (8)

$$S_z | S = 1/2, m_s = -1/2 \rangle = (-\hbar/2) | S = 1/2, m_s = -1/2 \rangle$$
 (9)

$$S_+|S=1/2, m_s=1/2\rangle = 0$$
 (10)

$$S_+|S=1/2, m_s=-1/2\rangle = \hbar|S=1/2, m_s=1/2\rangle$$
 (11)

$$S_{-}|S = 1/2, m_s = 1/2 \rangle = \hbar |S = 1/2, m_s = -1/2 \rangle$$
 (12)

$$S_{-}|S = 1/2, m_s = -1/2 \rangle = 0.$$
 (13)

[If you take PHYS 4443 you will derive these.] Similarly, for a spin-3/2 I the states obey:

$$I_z | I = 3/2, m_I = 3/2 \rangle = (3\hbar/2) | I = 3/2, m_I = 3/2 \rangle$$
 (14)

$$I_z | I = 3/2, m_I = 1/2 \rangle = (\hbar/2) | I = 3/2, m_I = 1/2 \rangle$$
 (15)

$$I_z | I = 3/2, m_I = -1/2 \rangle = (-\hbar/2) | I = 3/2, m_I = -1/2 \rangle$$
 (16)

$$I_z | I = 3/2, m_I = -3/2 \rangle = (-3\hbar/2) | I = 3/2, m_I = -3/2 \rangle$$
 (17)

$$I_{+}|I = 3/2, m_{I} = 3/2 \rangle = 0 \tag{18}$$

$$I_{+}|I = 3/2, m_{I} = 1/2 \rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = 3/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 3/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 1/2 \rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = 3/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 1/2 \rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = 3/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 1/2 \rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = 3/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 1/2 \rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = 3/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 1/2 \rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = 3/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 1/2 \rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = 3/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 3/2 \rangle = 1/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 3/2 \rangle = 1/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 3/2 \rangle = 1/2 \rangle = 1/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 3/2 \rangle = 1/2 \rangle = 1/2 \rangle = 1/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = 3/2 \rangle = 1/2 \rangle$$

$$I_{+}|I = 3/2, m_{I} = -1/2\rangle = \sqrt{7/2\hbar} |I = 3/2, m_{I} = 1/2\rangle$$

$$I_{+}|I = 3/2, m_{I} = -3/2\rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = -1/2\rangle$$
(20)
(21)

$$I = \frac{1}{2} - \frac{3}{2}, m_I = -\frac{3}{2}, m_I = -\frac{3}{2}, m_I = -\frac{3}{2}, m_I = -\frac{1}{2}, m_I =$$

$$I_{-}|I = 3/2, m_{I} = 5/2, m_{I} = 5/2, m_{I} = 5/2, m_{I} = 1/2, m_{I} = 1/2, m_{I} = 1/2, m_{I} = 1/2$$

$$I_{-}|I = 3/2, m_{I} = 1/2$$

$$(23)$$

$$I_{-}|I = 3/2, m_{I} = -1/2\rangle = \sqrt{3}\hbar |I = 3/2, m_{I} = -3/2\rangle$$
 (24)

$$I_{-}|I=3/2, m_{I}=-3/2\rangle = 0.$$
 (25)

[Again, you will derive this in PHYS 4443.] In recitation you had a chance to play with a computer algebra system which knows about these operators.

If we specify both the electronic and nuclear spin quantum numbers, we will write something like

$$|m_s = 1/2, m_I = 1/2\rangle = |S = 1/2, m_s = 1/2\rangle \otimes ||I = 3/2, m_I = 1/2\rangle.$$
 (26)

The S operators only act on the S quantum numbers, and similarly with I. For example

$$S_{-}|m_{s} = 1/2, m_{I} = 1/2 \rangle = \hbar |m_{s} = -1/2, m_{I} = 1/2 \rangle$$
(27)

or

$$I_{+}|m_{s} = 1/2, m_{I} = 1/2 \rangle = \sqrt{3}\hbar|m_{s} = 1/2, m_{I} = 3/2 \rangle.$$
(28)

Finally, you will find the following identity useful:

$$S \cdot I = S_z I_z + S_x I_x + S_y I_y \tag{29}$$

$$= S_z I_z + \frac{1}{2} \left( S_+ I_- + S_- I_+ \right) \tag{30}$$

**3.3.** There is a unique eigenstate of this Hamiltonian with  $m_f = m_s + m_I = 2$ , it is  $|m_s = 1/2, m_I = 3/2\rangle$ . Calculate its energy as a function of *B*. This linear behavior is similar to what happens if you just have a single level of a spin-1/2 electron – it is known as the Zeeman effect.

[Hint: Just act on this state with H, and use the given rules.]

**3.4.** Find the eigenstates and eigenvalues of this Hamiltonian which have  $m_f = 1$ . These are linear combinations of

$$|A\rangle = |m_s = 1/2, m_I = 1/2\rangle \tag{31}$$

$$|B\rangle = |m_s = -1/2, m_I = 3/2\rangle$$
 (32)

That is write  $\psi = \alpha |A\rangle + \beta |B\rangle$ , and write down a matrix eigenvalue equation for the vector

$$\vec{\psi} = \begin{pmatrix} \alpha \\ \beta \end{pmatrix}. \tag{33}$$

You will need to find the eigenvalues of that  $2 \times 2$  matrix.

Plot energy vs magnetic field Be sure to label the axes so we know what units you are using. Choose your scale so that the important features can be seen.

This is the Paschen-Bach effect, sometimes called the nonlinear Zeeman effect.

### Problem 4. Calculus of Variation

Some of you said that your analytical mechanics course did not formally define the math behind "Calculus of Variation". This problem works you through the formal definition. In a typical classical mechanics class one spends several weeks on this, so this is going to be a bit abbreviated.

Let y(x) be an arbitrary function. We can define a functional F[y] as an object which takes the function y and spits out a number. For example,

$$F[y] = \frac{1}{2} \int \frac{dx_1 \, dx_2}{|x_1 - x_2|} y(x_1) y(x_2). \tag{34}$$

The variational derivative is defined by

$$\frac{\delta F[y]}{\delta y(x)} = \lim_{\eta \to 0} \frac{F[y + \eta s_x] - F[y]}{\eta},\tag{35}$$

where  $s_x(t) = \delta(t - x)$  is a Dirac delta function. So in this example

$$F[y + \eta s_x] = \frac{1}{2} \int \frac{dx_1 dx_2}{|x_1 - x_2|} \left( y(x_1) + \eta \delta(x_1 - x) \right) \left( y(x_2) + \eta \delta(x_2 - x) \right).$$
(36)

**4.1.** Write  $(F[y + \eta s_x] - F[y])/\eta$  as an integral over  $x_1$  and  $x_2$ . Simplify the expression, but don't do the integrals. You should find two terms, one which is proportional to  $\eta$ , and one that is independent of  $\eta$ .

**4.2.** Since we will be taking the limit  $\eta \to 0$ , neglect the term proportional to  $\eta$ . Calculate the integrals for the other term, and hence evaluate

$$\frac{\delta F[y]}{\delta y(x)} = \lim_{\eta \to 0} \frac{F[y + \eta s_x] - F[y]}{\eta}.$$
(37)

4.3. Consider a more generic symmetric Kernel,

$$G[y] = \frac{1}{2} \int dt_1 \, dt_2 \, K(t_1, t_2) y(t_1) y(t_2), \tag{38}$$

where  $K(t_1, t_2) = K(t_2, t_1)$  is simply a function of two variables. Using any technique at your disposal, calculate

$$\frac{\delta G[y]}{\delta y(t)}.\tag{39}$$

Note: I changed the name of the dummy variables, but the logic should still be clear.

4.4. Consider a Kernel of the form:

$$K(t_1, t_2) = -m\delta''(t_1 - t_2) - m\omega^2\delta(t_1 - t_2)$$
(40)

where  $\delta(t)$  is the Dirac delta function, and

$$\delta''(t) = \frac{\partial^2}{\partial t^2} \delta(t). \tag{41}$$

By integrating by parts in the variables  $t_1$  and  $t_2$ , show that

$$G[y] = \int dt \left(\frac{1}{2}m(y'(t))^2 - \frac{1}{2}m\omega^2(y(t))^2\right)$$
(42)

4.5. By taking the result from problem 4.3, and substituting in Eq. (40), show that

$$\frac{\delta G[y]}{\delta y(t)} = -my''(t) - m\omega^2 y(t), \tag{43}$$

and hence Newton's equations for the simple harmonic oscillator can be written as

$$\frac{\delta G[y]}{\delta y(t)} = 0 \tag{44}$$

[Note, most classical mechanics courses use slightly different notation for this.]

## Problem 5. Feedback

5.1. How long did this homework take?

**5.2.** Which of the following words come to mind when you think about this homework (feel free to add your own words if you have something better): frustrating, fun, tedious, insightful, hard, easy, useful, useless, fair, unfair