

## P3317 HW from Lecture 20+21 and Recitation 12

Due Nov 20, 2018

### Problem 1. Landau Levels

Here we will find the energy eigenstates a quantum mechanical particle moving in two dimensions in a uniform magnetic field. That is we will consider the Schrodinger equation

$$E\psi = -\frac{\hbar^2}{2m} (\nabla - iq\mathbf{A})^2 \psi, \quad (1)$$

where  $\nabla \times \mathbf{A} = \mathbf{B}$ . We will take  $\mathbf{B}$  to point in the  $\hat{\mathbf{z}}$  direction, and be uniform in space. The vector potential is not unique, we can add a divergence of any function to it. This is often described as “gauge freedom.”

1.1. Show that

$$\mathbf{A} = \Lambda x \hat{\mathbf{y}}$$

, where  $\Lambda$  is a constant, corresponds to a uniform magnetic field. This is known as the *Landau gauge*, and it is particularly useful for this calculation.

1.2. Within the Landau gauge, the Hamiltonian involves  $\partial_x$ ,  $\partial_y$ , and  $x$ . It does not depend explicitly on  $y$ , so the momentum in the  $\hat{\mathbf{y}}$  direction is a good quantum number. Thus we write

$$\psi(x, y) = e^{ik_y y} \phi_{k_y}(x). \quad (2)$$

Plug this into Eq. (1) and get an equation for  $\phi_{k_y}(x)$ .

1.3. You should recognize that this equation is the same equation obeyed by the wave function for a 1D quantum harmonic oscillator. What is the frequency  $\omega_0$  of the oscillator? Thinking about the classical motion of a charged particle in a magnetic field, what is the physical significance of this frequency?

1.4. From your knowledge of the harmonic oscillator eigen-energies, what are the energies of all the single particle states? Note: since  $E$  does not depend on  $k_y$ , you have discovered a massive degeneracy.

### Problem 2. Guiding Centers

Consider a charged particle moving in 2D in a uniform magnetic field within the Landau Gauge:

$$H = \frac{1}{2m} [p_x^2 + (p_y - qBx)^2]. \quad (3)$$

Use Heisenberg equations of motion to calculate the rate of change of these four quantities:  $X = \langle \hat{x} \rangle$ ,  $Y = \langle \hat{y} \rangle$ ,  $\Pi_x = \langle p_x \rangle$ , and  $\Pi_y = \langle \hat{p}_y - qB\hat{x} \rangle$ . This should yield a set of four coupled differential equations which you can solve.

Hint 1: Start with the equations of motion for  $\Pi_x$  and  $\Pi_y$ . These will be closed. Solve them. Substitute the solutions into the equations for  $X$  and  $Y$ . Don't forget your constants of integration.

Hint 2: The Heisenberg equations of motion read:

$$\partial_t \langle \hat{O} \rangle(t) = \frac{1}{i\hbar} \int dr \psi^*(x, t) (\hat{O} \hat{H} - \hat{H} \hat{O}) \psi(x, t) \quad (4)$$

$$= \frac{1}{i\hbar} \langle [\hat{O}, \hat{H}] \rangle. \quad (5)$$

Hint 3: In case 2D confuses you – recall  $x$  commutes with  $y$  and  $p_y$ .

Hint 4: Recall the following:

$$[\hat{x}, \hat{p}] = i\hbar \quad (6)$$

$$[\hat{x}, \hat{p}^2] = 2i\hbar \hat{p} \quad (7)$$

$$[\hat{x}, \hat{x}] = 0 \quad (8)$$

$$[\hat{x}, \hat{x}^2] = 0 \quad (9)$$

$$[\hat{p}, \hat{x}] = -i \quad (10)$$

$$[\hat{p}, \hat{x}^2] = -2i\hbar \hat{x} \quad (11)$$

$$[\hat{p}, \hat{p}] = 0 \quad (12)$$

$$[\hat{p}, \hat{p}^2] = 0. \quad (13)$$

### Problem 3. Maxwell's Equations

We define the electromagnetic field tensor as

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu \quad (14)$$

where  $\mu, \nu = t, x, y, z$ , and  $A_t = -\phi$  is the scalar potential, and  $A_x, A_y, A_z$  are the components of the vector potential.

**3.1.** What are the following in terms of  $E$  and  $B$ :  $F_{xt}, F_{yt}, F_{zt}, F_{xy}, F_{yz}, F_{zx}$ ?

**3.2.** Consider the quantity

$$\Lambda_{xyz} = \partial_x F_{yz} + \partial_y F_{zx} + \partial_z F_{xy}. \quad (15)$$

By writing  $F$  in terms of  $A$ , show that  $\Lambda_{xyz} = 0$ . This should be part of one of Maxwell's equations. Which one is it?

**3.3.** Consider the quantity

$$\Lambda_{txy} = \partial_t F_{xy} + \partial_x F_{yt} + \partial_y F_{tx}. \quad (16)$$

By writing  $F$  in terms of  $A$ , show that  $\Lambda_{txy} = 0$ . This should be part of one of Maxwell's equations. You can complete that equation by looking at  $\Lambda_{tyz}$  and  $\Lambda_{tzy}$  (but you do not need to).

We will not prove it here, but it turns out that in this notation the other two Maxwell equations are

$$\partial_x F_{xt} + \partial_y F_{yt} + \partial_z F_{zt} = 0 \quad (17)$$

$$-\partial_t F_{tx} + \partial_y F_{yx} + \partial_z F_{zx} = 0 \quad (18)$$

$$-\partial_t F_{ty} + \partial_x F_{xy} + \partial_z F_{zy} = 0 \quad (19)$$

$$-\partial_t F_{tz} + \partial_x F_{xz} + \partial_y F_{yz} = 0 \quad (20)$$

$$(21)$$

#### Problem 4. Goldstone mode

Consider the following wave equation

$$\partial_t^2 \phi - \partial_x^2 \phi - m^2 \phi + \lambda^2 |\phi|^2 \phi = 0. \quad (22)$$

This is *not* a gauge theory.

**4.1.** Linearize Eq. (22) about  $\phi = \phi_0 = m/\lambda$ .

**4.2.** What are the spectra of excitations? You should find one massive (gapped) mode. You will also find a gapless mode – the Goldstone mode.

This is a generic result. Whenever you have a continuously degenerate ground state, you will have a gapless mode. For example, in a crystal there is a continuous degeneracy corresponding to where the first atom sits. Once you specify its location you know where all of them are. Thus a crystal should have a gapless mode. It in fact has 3 gapless modes: the two transverse acoustic phonons, and a single longitudinal acoustic phonon.

#### Problem 5. Feedback

**5.1.** How long did this homework take?

**5.2.** Which of the following words come to mind when you think about this homework (feel free to add your own words if you have something better): frustrating, fun, tedious, insightful, hard, easy, useful, useless, fair, unfair